

Performance and use of Mössbauer goodness-of-fit parameters: Response to spectra of varying signal/noise ratio and possible misinterpretations

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ABSTRACT

Simulated Mössbauer spectra of Lorentzian form have been fit with proper and improper models over a wide range of signal/noise (S/N) ratios in order to evaluate the characteristics of the goodness-of-fit parameters χ^2 and MISFIT as a function of sample concentration and run duration. Results demonstrate a marked difference in behavior. χ^2 may be reduced, suggesting an improved fit, if an incorrectly fit spectrum is degraded in quality. This is not the case with MISFIT. However, MISFIT may be defined in two slightly differing ways, each of which “diverges,” to large negative and positive values, respectively, at small S/N ratios. Failure to allow for these characteristics can lead to fundamental misinterpretations of an effect like “spectrometer drift” or the estimation of optimal data collection conditions.

The application of χ^2 , MISFIT, and Δ MISFIT parameters to fit evaluations in mineralogical spectroscopy is reviewed in light of the differing parameter responses.

INTRODUCTION

Recently Dyar (1984) has presented a comparison of Mössbauer spectroscopy facilities in an attempt to establish optimal run parameters and system-system disparities. Unfortunately, she failed to analyze the effects of varied signal/noise (S/N) ratio independently on the performance of the goodness-of-fit parameters themselves. As will be shown, changes in S/N can have large effects on the value of χ^2 or MISFIT for identical fit models. Hence her results in several cases require re-evaluation.

Dyar also neglected to utilize the Δ MISFIT parameter in the evaluation of her spectral fits. As has already been shown elegantly by Ruby (1973), and which will be briefly demonstrated in this paper, the great utility of the MISFIT parameter resides in comparison with its uncertainty, i.e., Δ MISFIT.

In the present work the goodness-of-fit parameters are defined and tested by application to the fitting of simulated Lorentzian Mössbauer spectra with correct and incorrect models.

GOODNESS-OF-FIT PARAMETERS

One of the crucial problems in evaluating any type of fit to spectral data is a proper answer to the question “How well is the model confirmed by the experimental data points?” This question follows from direct application of the scientific method, whereas the question “What model when applied to the data gives the best fit by some criterion?” is more generally the question that is answered.

The first question allows for testing of a (presumably) proper model, whereas the second simply looks for the best-fit model without testing its significance.

The general goodness-of-fit parameter that is calculated,

χ^2 , has difficulties when used with Mössbauer spectra since it is affected by the magnitude of the background count and, as will be demonstrated, can actually be reduced by degrading the quality of a spectrum to which an incorrect model is being fit. Hence χ^2 is not a good parameter for comparison of spectra collected at differing laboratories with differing experimental arrangements. Furthermore, as has been discussed by Law (1973), evaluation of χ^2 alone is insufficient to determine the appropriateness of the fitting model. The testing should always be done with independent confirming evidence.

To improve the situation, Ruby (1973) suggested a new goodness-of-fit parameter which he called “MISFIT.” This parameter appears to be much more insensitive to spectral quality than χ^2 . Moreover, comparison between the calculated MISFIT value and the uncertainty in the MISFIT (Δ MISFIT) can yield a clear answer to the first question raised above.

χ^2 is defined as

$$\chi^2 = \sum_{I=1}^N \frac{[Y_C(I) - Y_D(I)]^2}{Y_D(I)}, \quad (1)$$

where N is the number of data points, and the $Y_C(I)$ and the $Y_D(I)$ are the calculated and observed spectral values, respectively.

The reduced χ^2 (χ_R^2) is equal to χ^2 divided by the number of degrees of freedom of the least-squares fit, i.e., the number of data points minus the number of varied parameters, $N - n$:

$$\chi_R^2 = \chi^2 / (N - n). \quad (2)$$

Table 1. Results of least-squares fits to simulated Mössbauer spectra

Spectrum background count	Absorption	Type of fit model*	M (%)	ΔM (%)	M' (%)	$\Delta M'$ (%)	χ^2_R	S/N**
50 000	3.3%	100% L	-0.977	-1.031	13.622	16.074	1.00	7.38
200 000	3.3%	100% L	-0.246	-0.259	3.848	1.701	1.00	14.76
500 000	3.3%	100% L	-0.099	-0.104	1.583	0.436	1.00	23.34
1 000 000	3.3%	100% L	-0.049	-0.052	0.800	0.175	1.00	33.00
4 000 000	3.3%	100% L	-0.012	-0.013	0.202	0.036	1.00	66.00
10 000 000	3.3%	100% L	-0.005	-0.005	0.081	0.014	1.00	104.36
50 000	3.3%	85% L 15% G	-0.768	-1.085	14.028	16.774	1.01	7.38
200 000	3.3%	85% L 15% G	-0.110	-0.286	4.052	1.806	1.03	14.76
500 000	3.3%	85% L 15% G	0.014	0.123	1.726	0.472	1.07	23.34
1 000 000	3.3%	85% L 15% G	0.052	0.068	0.917	0.194	1.14	33.00
4 000 000	3.3%	85% L 15% G	0.075	0.024	0.294	0.044	1.45	66.00
10 000 000	3.3%	85% L 15% G	0.078	0.013	0.166	0.020	2.02	104.36
500 000	0.33%	100% L	-11.343	-10.715	59.055	273.51	1.00	2.33
500 000	1.0%	100% L	-1.135	-1.163	14.946	19.058	1.00	7.00
500 000	3.3%	100% L	-0.099	-0.104	1.583	0.436	1.00	23.34
500 000	10%	100% L	-0.010	-0.011	0.177	0.032	1.00	70.00
500 000	0.33%	85% L 15% G	-10.463	-11.079	59.928	279.66	1.00	2.33
500 000	1.0%	85% L 15% G	-0.915	-1.222	15.376	19.868	1.01	7.00
500 000	3.3%	85% L 15% G	0.014	0.123	1.726	0.472	1.07	23.34
500 000	10%	85% L 15% G	0.074	0.021	0.267	0.039	1.47	70.00

* L = Lorentzian, G = Gaussian.

** S/N = signal/noise ratio—see text for definition.

The MISFIT parameter is defined in terms of two variations on χ^2_R . The distance D is defined as

$$D = N(\chi^2_R - 1), \quad (3)$$

which, on substituting for χ^2_R becomes

$$D = \left\{ \frac{N}{N-n} \sum_{I=1}^N \frac{[Y_C(I) - Y_D(I)]^2}{Y_D(I)} \right\} - N. \quad (4)$$

The distance is the discrepancy or distance of the fit from the observed data.

Ruby (1973) neglected the size of n in formulating MISFIT so that D is simply

$$D = \sum_{I=1}^N \left\{ \frac{[Y_C(I) - Y_D(I)]^2}{Y_D(I)} - 1 \right\}. \quad (5)$$

The second variant on χ^2_R is the signal S , defined as the summation of the spectral data above (or below) some baseline. The definition of the baseline, Y_0 , is nontrivial, but neglecting this for the moment, the definition of S is

$$S = \sum_{I=1}^N \left\{ \frac{[Y_0 - Y_D(I)]^2}{Y_D(I)} - 1 \right\}. \quad (6)$$

MISFIT is the ratio of D/S from Equations 5 and 6.

The uncertainty in M is given by (Bevington, 1969):

$$\sigma_M = \left(\frac{\sigma_{D^2}}{D^2} + \frac{\sigma_{S^2}}{S^2} - 2 \frac{\sigma_{DS^2}}{DS} \right)^{1/2} M, \quad (7)$$

which is approximated by

$$\begin{aligned} \sigma_M &\approx \left[\frac{\sigma_{D^2}}{D^2} + \frac{\sigma_{S^2}}{S^2} \right]^{1/2} M \\ &\approx \left[\left(\frac{\Delta D}{D} \right)^2 + \left(\frac{\Delta S}{S} \right)^2 \right]^{1/2} M. \end{aligned} \quad (8)$$

Replacement of the ΔD and ΔS values yields (Ruby, 1973)

$$\Delta M = (1/S)[N(1 + M^2) + 4D(1 + M)]^{1/2}. \quad (9)$$

Another version of MISFIT that Ruby (1973, Eq. 7a) used for several examples is derived for cases of small absorption.

This formulation cannot produce negative values, is simpler to calculate, and has beneficial properties similar to those of the first MISFIT parameter. It can be written in the form

$$M' = \frac{\sum_{I=1}^N \frac{[Y_C(I) - Y_D(I)]^2}{Y_D(I)}}{\sum_{I=1}^N \frac{[Y_0 - Y_D(I)]^2}{Y_D(I)}}. \quad (10)$$

In order to display the differing characters of these goodness-of-fit parameters, synthetic Mössbauer spectra were generated with pure Lorentzian form and were fit with the known (generating) peak-shape functions and with incorrect functions. The spectra generation was carried out over a range of S/N that suitably mimicked the effect of varied run duration and sample concentration for typical Mössbauer experiments.

SPECTRUM SYNTHESIS

The synthetic Mössbauer spectra were generated by the Mössbauer fitting program MOSFT (Dollase, 1971, unpub.). The original set of data points comprise smooth Lorentzian curves with fixed background count and absorption fraction. The effect of varied S/N is added by scaling of the spectrum absorption, appropriate offsetting of the background count, and addition of random statistical noise

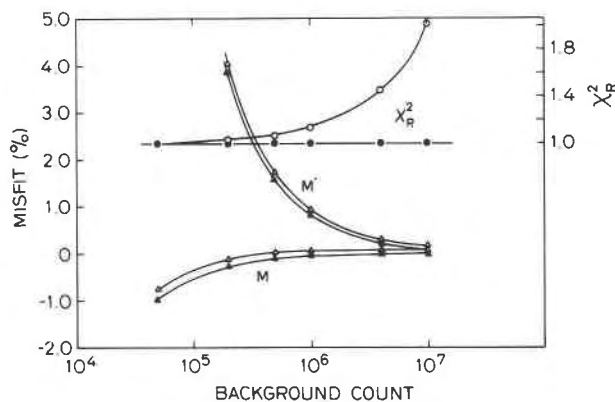


Fig. 1. MISFIT versus background count for fits of Lorentzian peaks (filled triangles) and 85% Lorentzian–15% Gaussian peaks (open triangles) to a Lorentzian simulation. χ_R^2 values for the same fits are plotted as filled and empty circles, respectively.

consistent with the magnitude of each resulting datum point according to the relation

$$A_{\text{vel}} + (A_{\text{vel}})^{1/2} R_{\text{norm}} = A'_{\text{vel}}$$

where

$$A_{\text{vel}} = (F_{\text{vel}} C_{\text{sf}}) - C_{\text{off}}$$

Here F_{vel} is the Lorentzian Mössbauer spectral line profile as generated, C_{sf} is a scaling factor to produce the desired absorption effect, C_{off} is a factor that offsets the background to be consistent with the absorption effect, A_{vel} is the scaled simulated Mössbauer spectrum, R_{norm} is a normally distributed random number with mean of 0 and standard deviation of 1, and A'_{vel} is the simulated Mössbauer spectrum with random statistical noise.

The use of a sufficiently random pseudorandom number generator was necessary to produce χ_R^2 values close to 1.00. The generator was tested rigorously to assure number sequences that closely approximated a normal distribution.

Synthetic spectra consisted of two equal Lorentzian lines summed to produce a doublet spectrum. Two sets of synthetic spectra were generated, one with varying background count and fixed sample absorption simulating the effect of run duration, and the second with fixed background count and varying sample absorption simulating the effect of sample concentration.

All spectra parameters are listed in Table 1.

SPECTRUM FITTING

Two types of fitting were performed. In the first type, all spectra were fit with 100% Lorentzian line shapes emulating the conditions of generation. Cycling during the refinement process was continued until there were at least two cycles without parameter change. Seven parameters were varied: position, area and width of each peak, and the background count.

In the second type, the spectra were fit with lines of 85% Lorentzian–15% Gaussian character, representing an incorrect model. The lineshape was produced by a sum

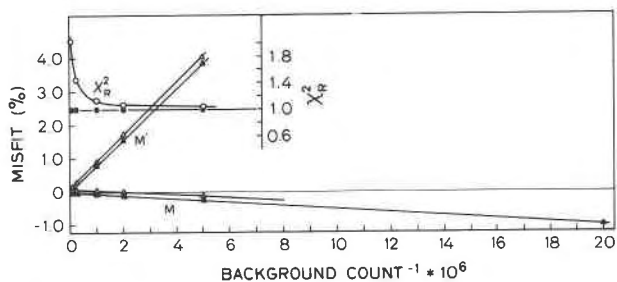


Fig. 2. MISFIT and χ_R^2 versus $1/\text{background count}$ for the fits of the simulations. Symbols as in Fig. 1.

of Lorentzian and Gaussian functions often referred to as a pseudo-Voigt profile (Hecq, 1981). As in the first type, the same seven parameters were varied, and the same requirement on fit convergence was used to determine the number of refinement cycles.

All of the goodness-of-fit parameters were calculated for each refined fit, as well as the standard deviations of these and all varied parameters.

RESULTS

The goodness-of-fit parameters resulting from the two types of fits to the first set of synthetic spectra are plotted in Figures 1 and 2 as a function of background count and the reciprocal of the background count, respectively.

Considering the Lorentzian-to-Lorentzian fit first, the value of χ_R^2 is found to be constant at 1.00 as should be expected for a random background and perfectly fit data. The MISFIT parameter M is found to increase from a negative value of -0.977% at the 50 000-count background level to -0.0049% at 10 000 000 background counts. M' , in contrast, decreases from 13.62 to 0.014% over the same range. Extrapolation to infinite background counts, in Figure 2, indicates that both MISFIT parameters would be zero in this limit.

In the case of the 85% Lorentzian–15% Gaussian fit to the Lorentzian spectra, the χ_R^2 values increase with improvement in the S/N ratio at greater background counts. The MISFIT parameter M increases with increasing background count, but rises to finite positive values instead of extrapolating to zero at infinite background count. M' decreases with increasing background count and extrapolates to the same finite positive value as M with infinite background count.

The goodness-of-fit parameters resulting from the two types of fits to the second set of synthetic spectra are plotted in Figure 3 as a function of percent absorption at fixed (500 000) background count. The characteristic behavior of each of the parameters is similar to that seen in Figure 1, except that the larger range of S/N explored produces wider parameter variation. For comparison, the complete set of results is plotted in Figure 4 as a function of S/N. Note that for this plot the signal is defined as the number of counts below background for the maximum spectral absorption, which is the usual spectral definition of signal, and the noise is the magnitude of the square

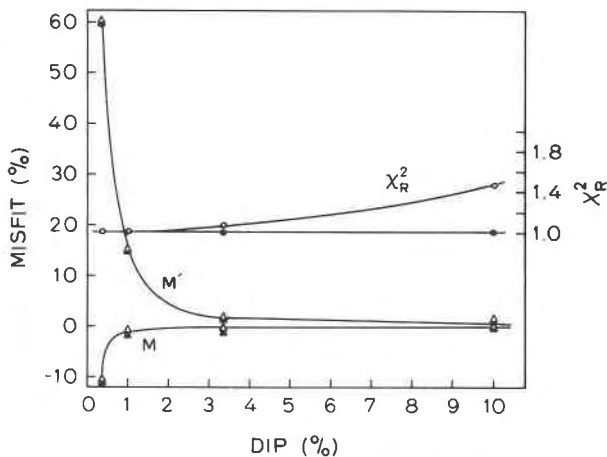


Fig. 3. MISFIT and χ^2_R versus Mössbauer effect absorption for fits to simulated Mössbauer spectra. Symbols as in Fig. 1.

root of the background count. The S/N ratio appears to be the appropriate variable for analyzing the behavior of Mössbauer goodness-of-fit parameters.

DISCUSSION

The importance of having a goodness-of-fit parameter other than χ^2 is evidenced by the behavior of χ^2 when an incorrect fitting model is used. In this case, χ^2 will continuously increase the longer a spectrum is collected. Hence, a poor fitting model may appear better than another superior model if the data are sufficiently poor. Since fitting models for mineralogical spectra are almost always at least as incorrect as those used in the present simulation because of the usual neglect of sample site variations, strain, blackness effects, cosine error, and so on, the trend in χ^2 seen for the synthetic spectra is generally applicable. Thus χ^2 values must be used cautiously for model testing.

On the other hand, if several runs of varying duration can be performed, the rate of change in χ^2 with background count affords some estimate of model validity.

The MISFIT parameters are less sensitive to S/N effects over the usual range of Mössbauer experiments than χ^2 , but are markedly sensitive at low S/N ratios. In particular, the negative values produced are artifacts of the formulation of MISFIT and do not represent improved fits. However, the extrapolation of MISFIT values to infinite signal/noise (or background counts) produces an absolute measure of fit quality.

The negative values can lead to misinterpretations. Dyar (1984) collected spectral data as a function of run duration which display an identical trend in M as the incorrect fit model presented in this work. Dyar attributed the effects of increasing M with run duration to spectrometer drift, but the full magnitude of the MISFIT change she observed is consistent with the dependence of M on S/N alone. Thus variation in M cannot be used to test for spectrometer drift if effects on spectral quality are small.

Dyar (1984) also concluded that the optimum run du-

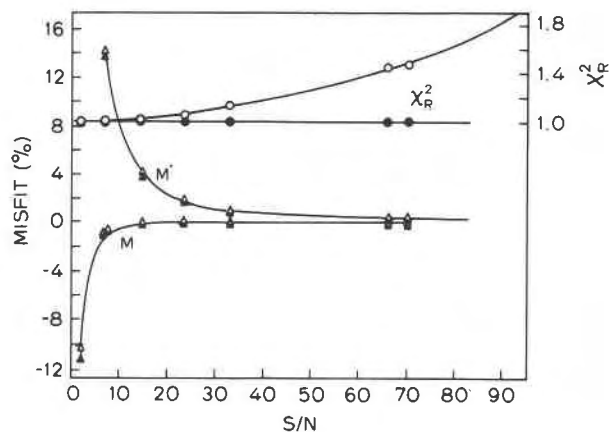


Fig. 4. MISFIT and χ^2_R versus the signal/noise ratio in fits to simulated Mössbauer spectra. Symbols as in Fig. 1.

ration for Mössbauer spectra occurred for M near zero. This conclusion is based on the false assumption that M is approximately zero for well-fit high-quality data. It is, in fact, true only for perfect fits in the limit of infinite S/N, both of which are not achieved in practice.

What is true is that spectrum quality must constantly improve, though at a slower and slower rate, with run duration. If spectrometer drift occurs, this is ordinarily manifested by alteration of spectral line shapes. Line-shape changes can be specified by fitting procedures or by Fourier deconvolution methods.

Another methodology consideration treated by Dyar (1984) is that of the optimum sample concentration. This was evaluated using the zero point in plots of MISFIT versus sample concentration for fixed run durations. For the reasons already outlined, this method should not be used to determine optimum experimental conditions. The value obtained by Dyar (1984) for a silicate sample is fortuitously close to theoretical estimates of ideal sample concentration, whereas her value for a sulfide sample is much too small.

USE OF ΔM AND THE Y_0 PROBLEM

Although the foregoing discussion indicates anomalies in the character of M and M' , these are not severe enough to obviate their virtues. The real utility in using MISFIT lies in the comparison of MISFIT and Δ MISFIT. Since Δ MISFIT indicates the uncertainty in M (and this is due mainly to the quality of the data) while M is generally much less sensitive to data quality, it reveals the degree to which the fit model is tested. For example, in the 85% Lorentzian-15% Gaussian fits to the synthetic spectra the value of M and ΔM for the 500 000-count 3.3% absorption spectrum fit are 0.014 and 0.123%, respectively. An M of 0.014% represents an excellent fit, but a ΔM of 0.123% indicates that the noise in the data would permit of M values up to 0.37% larger (3 standard deviations). Hence the model is not severely tested, and the low M value is, by itself, deceptive. In the case of the analogous 10 000 000-

count spectrum, M and ΔM are 0.078% and 0.013%, respectively. The M value is now "on target" and indicates a good fit with a small but definite flaw. The ΔM value indicates high-quality data and, since it is considerably smaller than M , good model testing.

This leads to two requirements for testing of any fitting models: (1) minimization of χ^2 and MISFIT and (2) ΔM much smaller than M .

A model that reduces χ^2 or M by 50% over some other model (on the same data) is only really superior to the other if ΔM is significantly smaller than 50% of M . Ruby (1973) provided several examples of the use of M and ΔM in deciphering fit and data quality.

A problem arises in calculating MISFIT accurately due to the difficulty in obtaining proper Y_0 values. Y_0 is known exactly for synthetic spectra, but if refined in a least-squares procedure on experimental data, the Y_0 value obtained is dependent on the part of the background that is refined. For example, if a sharp absorption occurs on a broad flat background and only data points far from the absorption are used to refine Y_0 , then the Y_0 value obtained is a good representation of the true value. Since this is difficult in practice, Y_0 is usually refined with data points incorporating some of the absorption "tail." This tends to shift Y_0 slightly toward the top of the absorption feature. In Mössbauer spectra Y_0 is thus reduced and S is smaller than the true value. Further, the shape of the absorption tail will obviously affect the refined Y_0 value. Hence Y_0 and any peak shape parameters will be correlated.

These concerns illustrate part of the difficulty in defining and utilizing a general goodness-of-fit parameter. Despite their weaknesses, however, it is clear that MISFIT and Δ MISFIT provide a useful alternative to χ^2 . The safest application of these parameters appears to be in combination where possible.

CONCLUSIONS

(1) The general goodness-of-fit parameter, χ^2 , has reduced value in judging fit quality since it is strongly dependent on data quality if an improper fit model is being utilized. (2) MISFIT is less sensitive than χ^2 to data quality but not insensitive to it. It has anomalous behavior for fits to low-quality data that must be taken into account when using it alone for judging fit perfection. (3) M comparison with ΔM is of high value in testing fit models. Use of M or χ^2 alone are most useful if spectra from the same sample are available with varied S/N ratios. (4) Attention should be paid to the proper refinement of Y_0 in order to estimate M , ΔM , and S correctly.

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REFERENCES

- Bevington, P.R. (1969) Data reduction and error analysis for the physical sciences. McGraw-Hill, New York.
- Dyar, M.D. (1984) Precision and interlaboratory reproducibility of measurements of the Mössbauer effect in minerals. *American Mineralogist*, 69, 1127–1144.
- Hecq, M. (1981) A fitting method for X-ray diffraction profiles. *Journal of Applied Crystallography*, 14, 60–61.
- Law, A.D. (1973) Critical evaluation of statistical "best fits" to Mössbauer spectra. *American Mineralogist*, 58, 128–131.
- Ruby, S.L. (1973) Why misfit when you already have χ^2 ? In T.J. Gruverman and C.W. Seidel, Eds. *Mössbauer effect methodology*, 8, 263–276. Plenum Press, New York.

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